

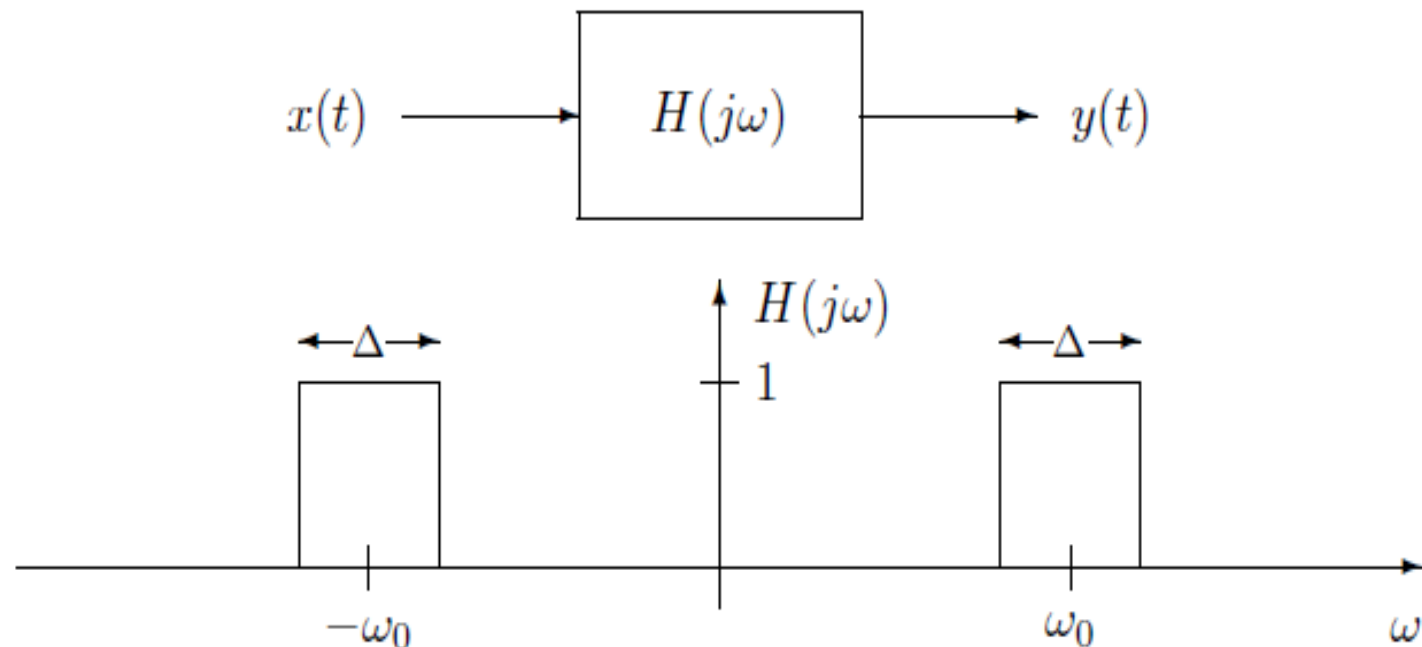
Signal power and power spectral  
density,  
properties of power spectral  
density

Motivated by situations in which  $x(t)$  is the voltage across (or current through) a unit resistor, we refer to  $x^2(t)$  as the *instantaneous power* in the signal  $x(t)$ . When  $x(t)$  is WSS, the *expected* instantaneous power is given by

$$E[x^2(t)] = R_{xx}(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} S_{xx}(j\omega) d\omega ,$$

where  $S_{xx}(j\omega)$  is the CTFT of the autocorrelation function  $R_{xx}(\tau)$ . Furthermore, when  $x(t)$  is ergodic in correlation, so that time averages and ensemble averages are equal in correlation computations, then (10.1) also represents the time-average power in any ensemble member. Note that since  $R_{xx}(\tau) = R_{xx}(-\tau)$ , we know  $S_{xx}(j\omega)$  is always *real* and *even* in  $\omega$ ; a simpler notation such as  $P_{xx}(\omega)$  might therefore have been more appropriate for it, but we shall stick to  $S_{xx}(j\omega)$  to avoid a proliferation of notational conventions, and to keep apparent the fact that this quantity is the Fourier transform of  $R_{xx}(\tau)$ .

The integral above suggests that we might be able to consider the expected (instantaneous) power (or, assuming the process is ergodic, the time-average power) in a frequency band of width  $d\omega$  to be given by  $(1/2\pi)S_{xx}(j\omega) d\omega$ . To examine this thought further, consider extracting a band of frequency components of  $x(t)$  by passing  $x(t)$  through an ideal bandpass filter, shown in Figure



Because of the way we are obtaining  $y(t)$  from  $x(t)$ , the expected power in the output  $y(t)$  can be interpreted as the expected power that  $x(t)$  has in the selected passband. Using the fact that

$$S_{yy}(j\omega) = |H(j\omega)|^2 S_{xx}(j\omega) ,$$

we see that this expected power can be computed as

$$E\{y^2(t)\} = R_{yy}(0) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} S_{yy}(j\omega) d\omega = \frac{1}{2\pi} \int_{\text{passband}} S_{xx}(j\omega) d\omega .$$

Thus

$$\frac{1}{2\pi} \int_{\text{passband}} S_{xx}(j\omega) d\omega$$

is indeed the expected power of  $x(t)$  in the passband. It is therefore reasonable to call  $S_{xx}(j\omega)$  the **power spectral density (PSD)** of  $x(t)$ . Note that the instantaneous power of  $y(t)$ , and hence the expected instantaneous power  $E[y^2(t)]$ , is always nonnegative, no matter how narrow the passband. It follows that, in addition to being real and even in  $\omega$ , the PSD is always nonnegative,  $S_{xx}(j\omega) \geq 0$  for all  $\omega$ . While the PSD  $S_{xx}(j\omega)$  is the Fourier transform of the autocorrelation function, it

is useful to have a name for the *Laplace* transform of the autocorrelation function; we shall refer to  $S_{xx}(s)$  as the *complex* PSD.

Exactly parallel results apply for the DT case, leading to the conclusion that  $S_{xx}(e^{j\Omega})$  is the power spectral density of  $x[n]$ .